

Particles which are Identical, Indistinguishable, Integral or zero spin - (Boson)

Examples: photons, phonons, α -particles, They do not obey

Pauli's Exclusion Principle.

Consider system of Bosons in thermal equilibrium at temperature 'T'. Let U be the energy of the system and N the number of particles in system. both U and N are constant.

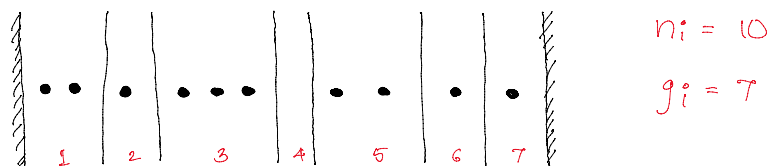
Let E_1, E_2, \dots, E_n be the energies of different particle states in increasing order and let g_1, g_2, \dots be their degeneracies respectively. Let n_i be the particles in energy state E_i and so on. so the conservation of particles demands that.

$$n_1 + n_2 + n_3 + \dots = \sum n_i = N = \text{constant}.$$

The conservation of energy says

$$\sum n_i E_i = U = \text{constant}.$$

Number of ways in which n_i particles may be distributed in energy state E_i which has degeneracy ' g_i '.



Note that $(g_i - 1)$ objects are required to split n_i particles into g_i compartments.

Therefore there are $(n_i + g_i - 1)!$ possible permutations that we can make within $n_i + g_i - 1$ objects, but all these permutations are not distinct because $n_i!$ permutations of n_i particles among themselves and $(g_i - 1)!$ permutations of $g_i - 1$ partitions among themselves do not affect distribution and they are irrelevant.

\therefore The number of distinct ways of arranging n_i bosons in g_i quantum state is given by

$$\frac{(n_i + g_i - 1)!}{n_i! (g_i - 1)!}$$

It holds for all values of 'i' so the total number of microstates accessible to the system for this distribution is given by.

$$\Omega = \prod_i \frac{(n_i + g_i - 1)!}{n_i! (g_i - 1)!}$$

\prod stands for product of function over 'i'

\prod_i stands for product of function over 'i'

\therefore Probability of distribution

$$W = \frac{\Omega}{\Omega_{\text{total}}}$$

Ω_{total} is the sum of accessible microstates considering all distributions possible for the system.

Most probable distribution:

$$\ln \Omega = \sum_i \ln(n_i + g_i - 1)! - \sum_i \ln n_i! - \sum_i \ln(g_i - 1)!$$

Using Stirling's approximation. ($\ln x! = x \ln x - x$)

$$\begin{aligned} \ln \Omega = & \sum_i (n_i + g_i - 1) \ln(n_i + g_i - 1) - \sum_i (n_i + g_i - 1) - \sum_i n_i \ln n_i + \sum_i n_i \\ & - \sum_i (g_i - 1) \ln(g_i - 1) + \sum_i (g_i - 1) \end{aligned}$$

$$\begin{aligned} \delta \ln \Omega = & \sum_i \ln(n_i + g_i - 1) \delta n_i + \sum_i (n_i + g_i - 1) \delta \ln(n_i + g_i - 1) \\ & - \sum_i n_i \delta \ln n_i - \sum_i \ln n_i \delta n_i \end{aligned}$$

Since $\delta \ln x = \frac{1}{x} \delta x$, we have.

$$\therefore \delta \ln \Omega = \sum_i \ln(n_i + g_i - 1) \delta n_i + \sum_i \delta n_i - \sum_i \delta n_i - \sum_i \ln n_i \delta n_i$$

Neglecting 1 in comparison to $n_i + g_i$ as $n_i + g_i \gg 1$

$$\delta \ln \Omega = \sum_i \ln \frac{n_i + g_i}{n_i} \delta n_i$$

\therefore for most probable distribution $\delta \ln \Omega = 0$

$$\therefore \sum_i \ln \frac{n_i + g_i}{n_i} \delta n_i = 0$$

All δn_i are not independent.

$$\therefore \sum_i \delta n_i = 0 \quad \text{conservation of particles}$$

$$\therefore \sum_i E_i \delta n_i = 0 \quad \text{conservation of energy}$$

These equations have to be satisfied simultaneously which can be done by Lagrangian method of undetermined multipliers

$$\therefore \sum_i \left(\ln \frac{n_i + g_i}{n_i} - \alpha - \beta E_i \right) \delta n_i = 0$$

$$\therefore \ln \frac{n_i + g_i}{n_i} - \alpha - \beta E_i = 0 \quad \text{for each } i$$

$$\therefore \ln \frac{n_i + g_i}{n_i} = \alpha + \beta E_i$$

$$\therefore \frac{n_i + g_i}{n_i} = e^{\alpha + \beta E_i}$$

$$\therefore 1 + \frac{g_i}{n_i} = e^{\alpha} \cdot e^{\beta E_i}$$

$$\therefore \frac{g_i}{n_i} = e^{\alpha} e^{\beta E_i} - 1$$

$$\therefore n_i = \frac{g_i}{e^{\alpha} e^{\beta E_i} - 1} \quad \left(\beta = \frac{1}{kT} \right)$$

\therefore This is BE distribution law.

$$\therefore n(\epsilon) d\epsilon = \frac{g(\epsilon) d\epsilon}{e^{\alpha} e^{\beta \epsilon} - 1}$$

where $n(\epsilon) d\epsilon$ is the number bosons having energy between ϵ & $\epsilon + d\epsilon$.